

# Differential expansion of space and the Hubble flow anisotropy

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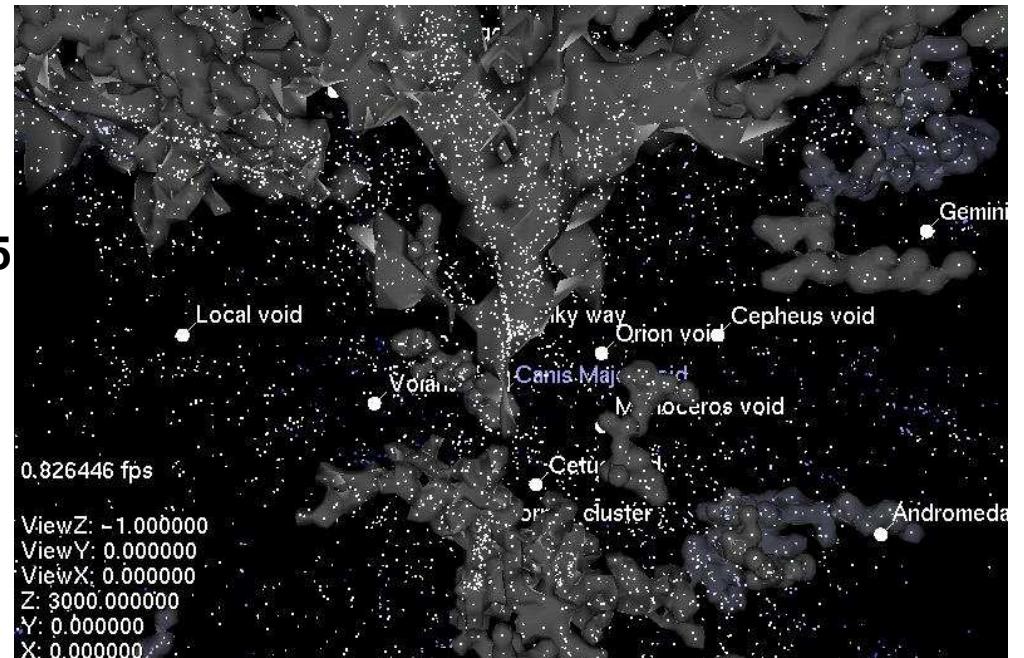
DLW, P R Smale, T Mattsson and R Watkins

**Phys. Rev. D 88 (2013) 083529**

J H McKay & DLW: **MNRAS 457 (2016) 3285**

K Bolejko, M A Nazer & DLW:

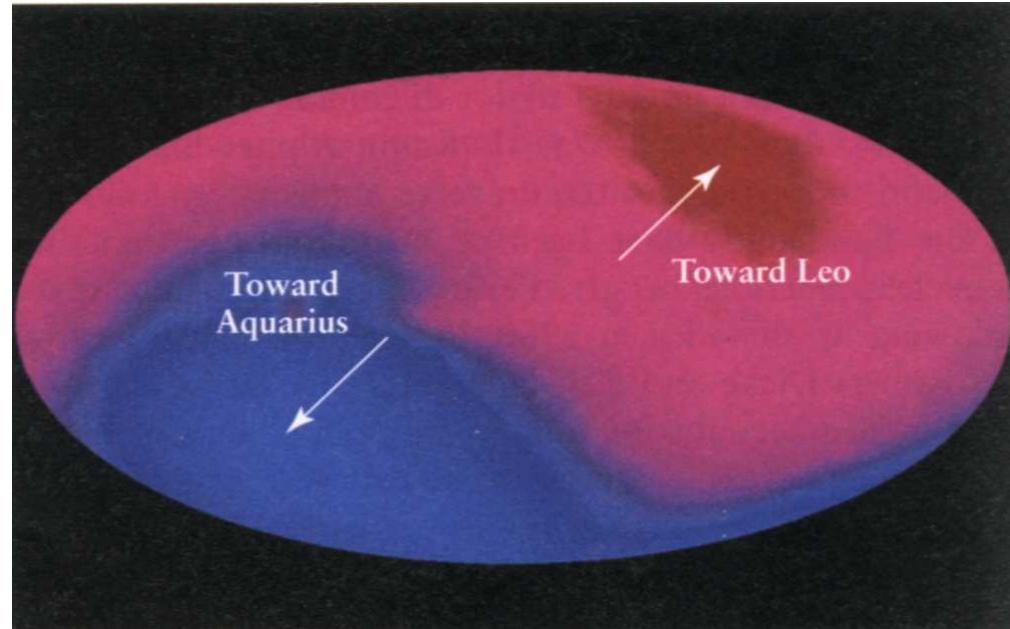
**JCAP 06 (2016) 035**



# Outline of talk

- Test of FLRW on  $\lesssim 100 h^{-1} \text{Mpc}$  scales  
*Result: alternative to large bulk flows - GR differential expansion on  $\lesssim 65 h^{-1} \text{Mpc}$  scales, supported by simulations, with proposal to test impact on large angle CMB anomalies*
- Use exact inhomogeneous solutions of Einstein's equations for structures on  $\lesssim 70 h^{-1} \text{Mpc}$  scales; asymptotic to Planck normalized FLRW model on larger scales: *effective* model for large scale light propagation
- Use Szekeres model: most general dust solution, reduces to spherically symmetric inhomogeneity (Lemaître–Tolman–Bondi (LTB) model) in a limit
- Trace rays from CMB and mock COMPOSITE catalogues for 4534 galaxy clusters

# Cosmic Microwave Background dipole



- Special Relativity: motion in a thermal bath of photons

$$T' = \frac{T_0}{\gamma(1 - (v/c) \cos \theta')}$$

- 3.37 mK dipole:  $v_{\text{Sun-CMB}} = 371 \text{ km s}^{-1}$  to  $(264.14^\circ, 48.26^\circ)$ ; splits as  $v_{\text{Sun-LG}} = 318.6 \text{ km s}^{-1}$  to  $(106^\circ, -6^\circ)$  and  $v_{\text{LG-CMB}} = 635 \pm 38 \text{ km s}^{-1}$  to  $(276.4^\circ, 29.3^\circ) \pm 3.2^\circ$

# Peculiar velocity formalism

- Standard framework, FLRW + Newtonian perturbations, assumes peculiar velocity field

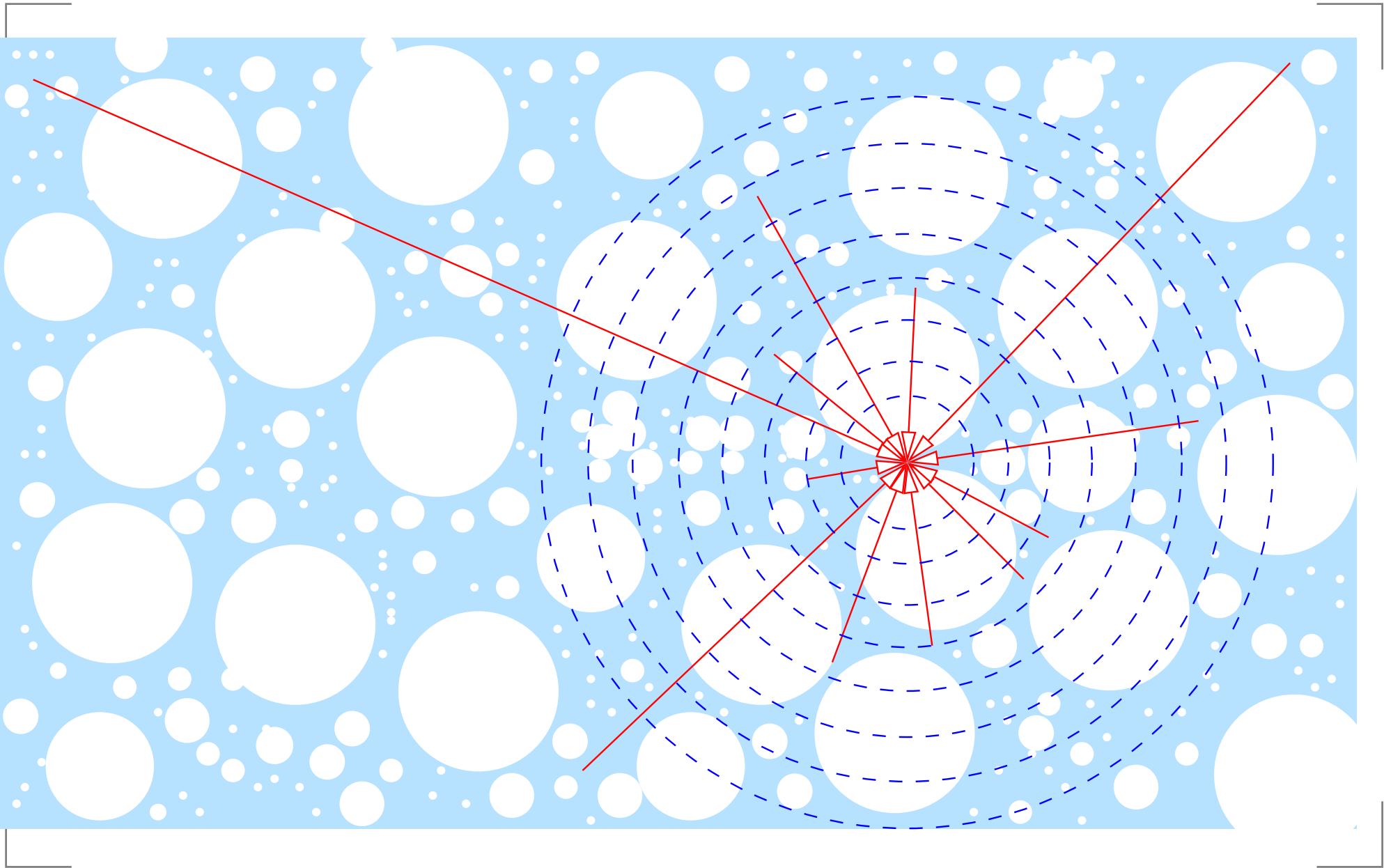
$$v_{\text{pec}} = cz - H_0 r$$

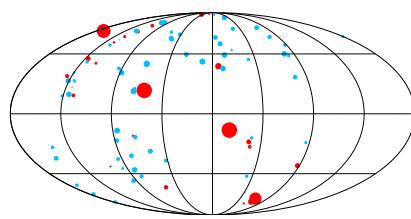
generated by

$$\mathbf{v}(\mathbf{r}) = \frac{H_0 \Omega_{M0}^{0.55}}{4\pi} \int d^3\mathbf{r}' \delta_m(\mathbf{r}') \frac{(\mathbf{r}' - \mathbf{r})}{|\mathbf{r}' - \mathbf{r}|^3}$$

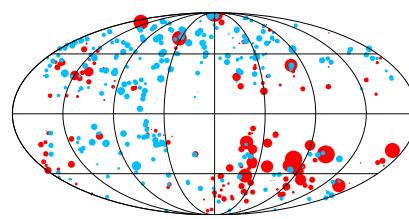
- 3 decades of debate on convergence of  $\mathbf{v}(\mathbf{r})$  to velocity of LG w.r.t. CMB frame; Direction agreed, not amplitude or scale (Lavaux et al 2010; Bilicki et al 2011...)
- The debate continues: Hess & Kitaura arXiv:1412.7310, Springob et al arXiv:1511.04849, ...

# Apparent Hubble flow variation

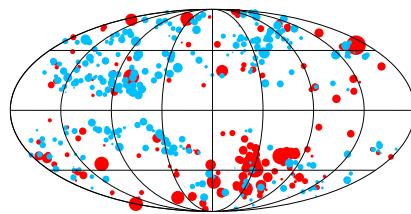




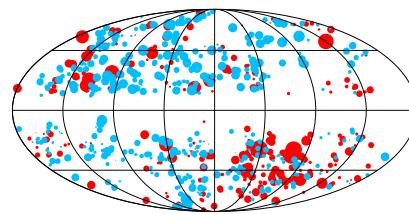
(a) 1:  $0 - 12.5 h^{-1}$  Mpc  $N = 92$ .



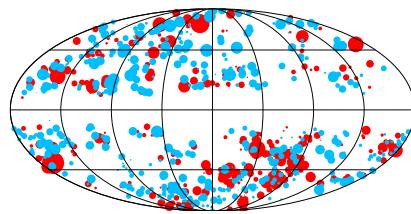
(b) 2:  $12.5 - 25 h^{-1}$  Mpc  $N = 505$ .



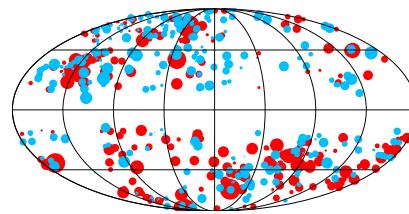
(c) 3:  $25 - 37.5 h^{-1}$  Mpc  $N = 514$ .



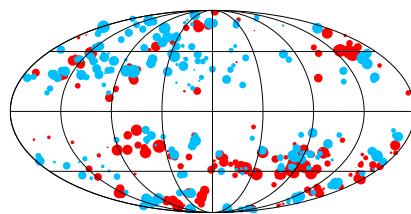
(d) 4:  $37.5 - 50 h^{-1}$  Mpc  $N = 731$ .



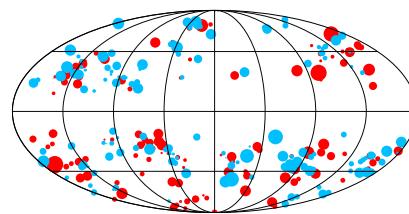
(e) 5:  $50 - 62.5 h^{-1}$  Mpc  $N = 819$ .



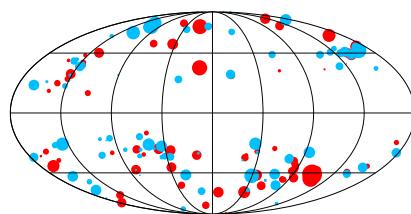
(f) 6:  $62.5 - 75 h^{-1}$  Mpc  $N = 562$ .



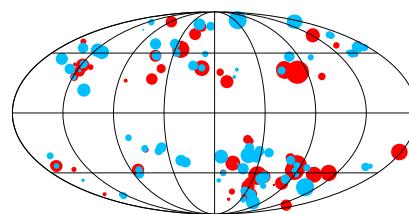
(g) 7:  $75 - 87.5 h^{-1}$  Mpc  $N = 414$ .



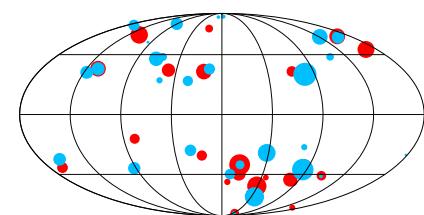
(h) 8:  $87.5 - 100 h^{-1}$  Mpc  $N = 304$ .



(i) 9:  $100 - 112.5 h^{-1}$  Mpc  $N = 222$ .

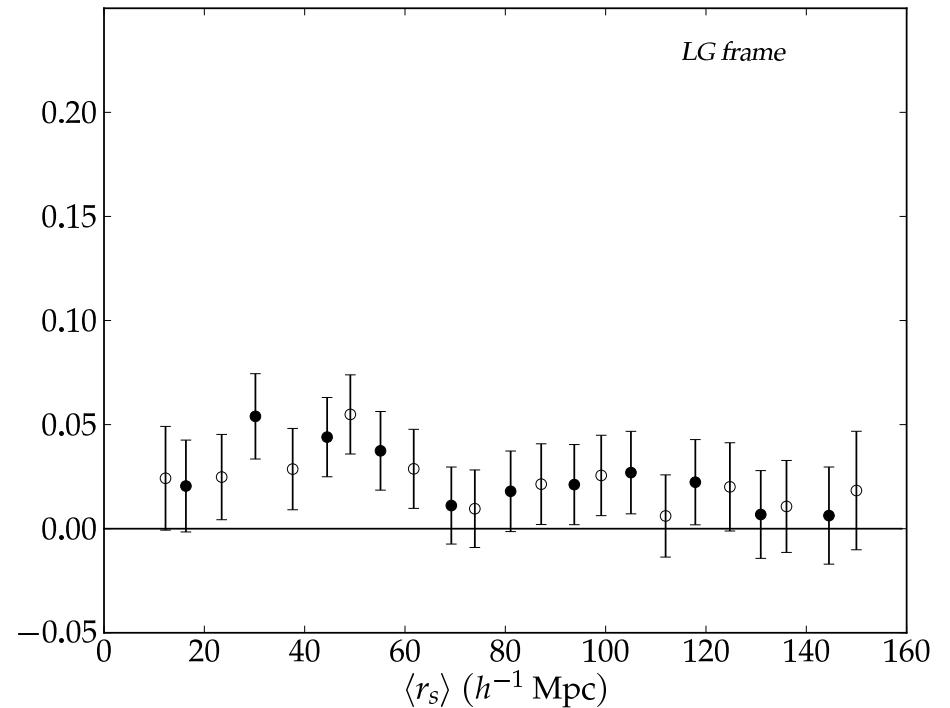
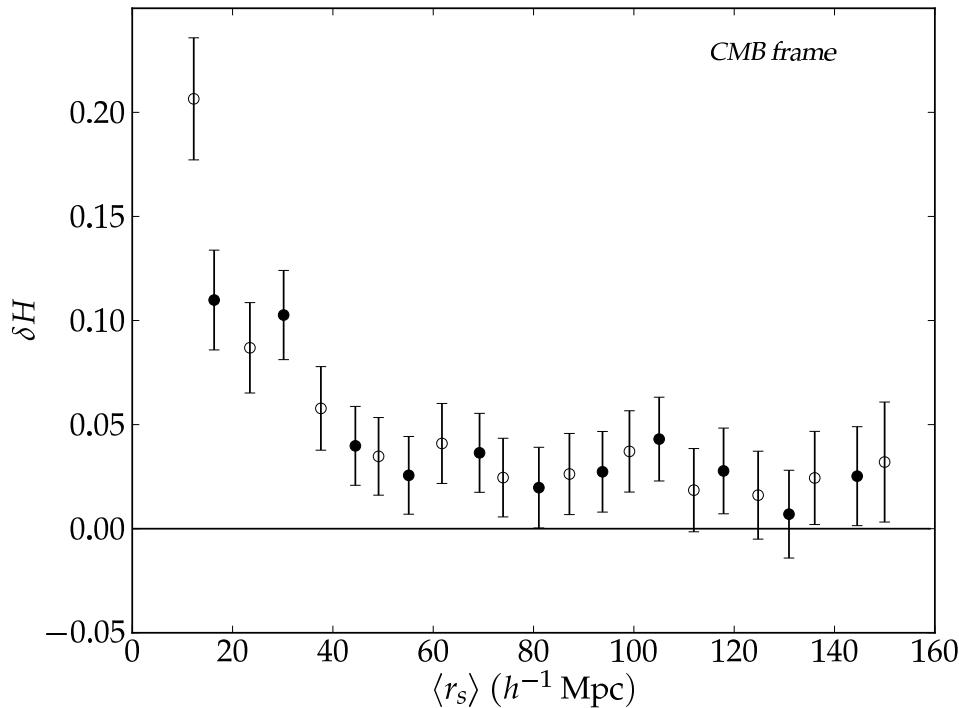


(j) 10:  $112.5 - 156.25 h^{-1}$  Mpc  $N = 280$ .



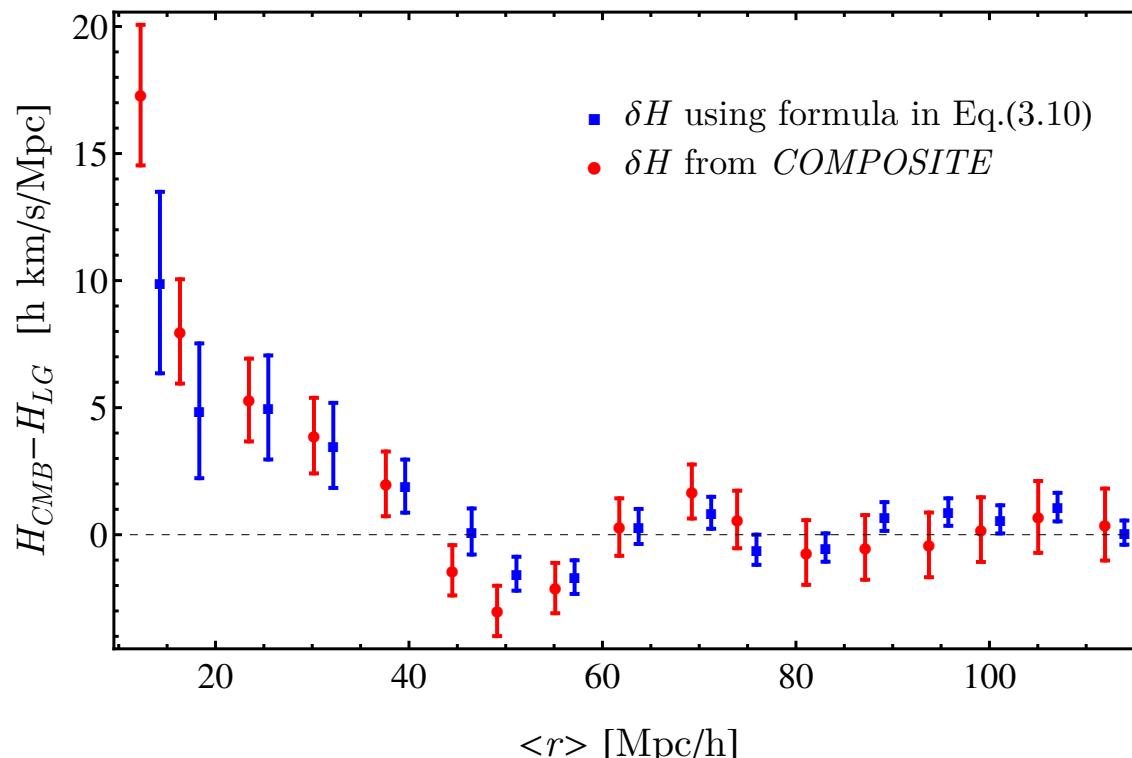
(k) 11:  $156.25 - 417.4 h^{-1}$  Mpc  $N = 91$ .

# Radial variation $\delta H_s = (H_s - H_0)/H_0$



- Two choices of shell boundaries (closed and open circles); for each choice data points uncorrelated
- Result: Hubble expansion is very significantly more uniform in LG frame than in CMB frame:  $\ln B > 5$ ; (except for  $40 \lesssim r \lesssim 60 h^{-1} \text{Mpc}$ ).

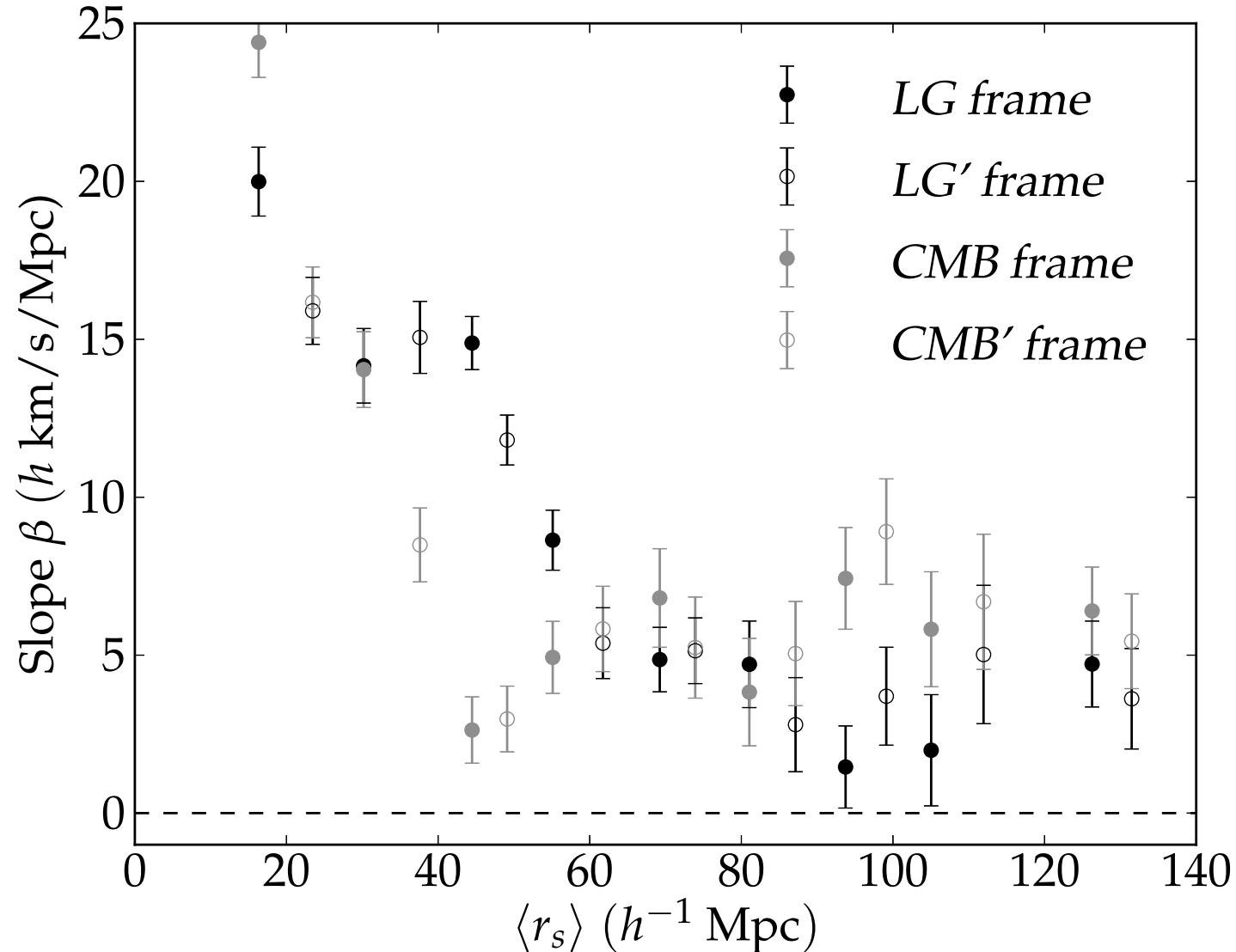
# Boost offset and deviation



- Kraljic and Sarkar (2016, to appear). FLRW + Newtonian N-body simulation with bulk flow  $\mathbf{v}_{\text{bulk}}(r)$

$$H'_s - H_s \sim \frac{|\mathbf{v}|^2 - 2\mathbf{v} \cdot \mathbf{v}_{\text{bulk}}(r)}{3H_0 \langle r^2 \rangle}$$

# Value of $\beta$ in $\frac{cz}{r} = H_0 + \beta \cos \phi$



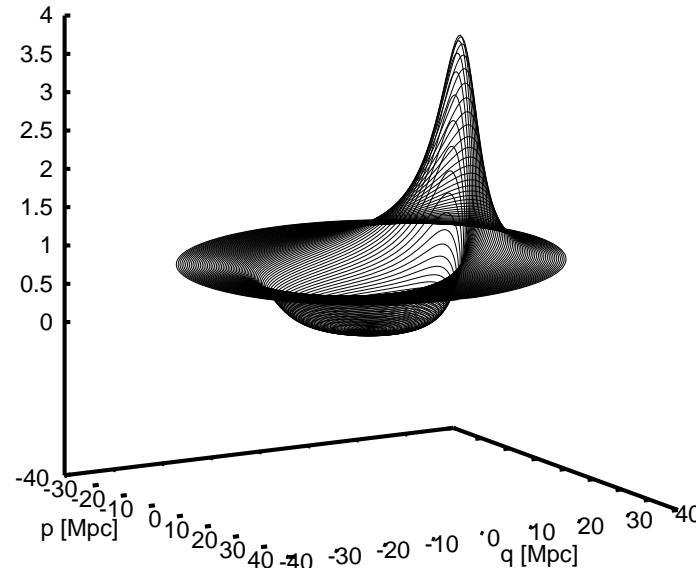
# Ray tracing: Szekeres model (1975)

$$\begin{aligned}
 ds^2 &= c^2 dt^2 - \frac{\left(R' - R \frac{\mathcal{E}'}{\mathcal{E}}\right)^2}{1-k} dr^2 - \frac{R^2}{\mathcal{E}^2} (dp^2 + dq^2), \\
 \mathcal{E}(r, p, q) &= \frac{1}{2S}(p^2 + q^2) - \frac{P}{S}p - \frac{Q}{S}q + \frac{P^2}{2S} + \frac{Q^2}{2S} + \frac{S}{2}, \\
 \dot{R}^2 &= -k(r) + \frac{2M(r)}{R} + \frac{1}{3}\Lambda R^2, \\
 t - t_B(r) &= \int_0^R d\tilde{R} \left[ -k + 2M/\tilde{R} + \frac{1}{3}\Lambda \tilde{R}^2 \right]^{-1/2}, \\
 \kappa\rho &= \frac{2(M' - 3M\mathcal{E}'/\mathcal{E})}{R^2(R' - R\mathcal{E}'/\mathcal{E})}.
 \end{aligned}$$

where  $' \equiv \partial/\partial r$ ,  $\cdot \equiv \partial/\partial t$ ,  $R = R(t, r)$ ,  $k = k(r) \leq 1$ ,  $S = S(r)$ ,  $P = P(r)$ ,  $Q = Q(r)$ ,  $M = M(r)$ . Above eqns satisfied but functions are otherwise arbitrary. We take  $t_B(r) = 0$ .

# Ray tracing: Szekeres model

density

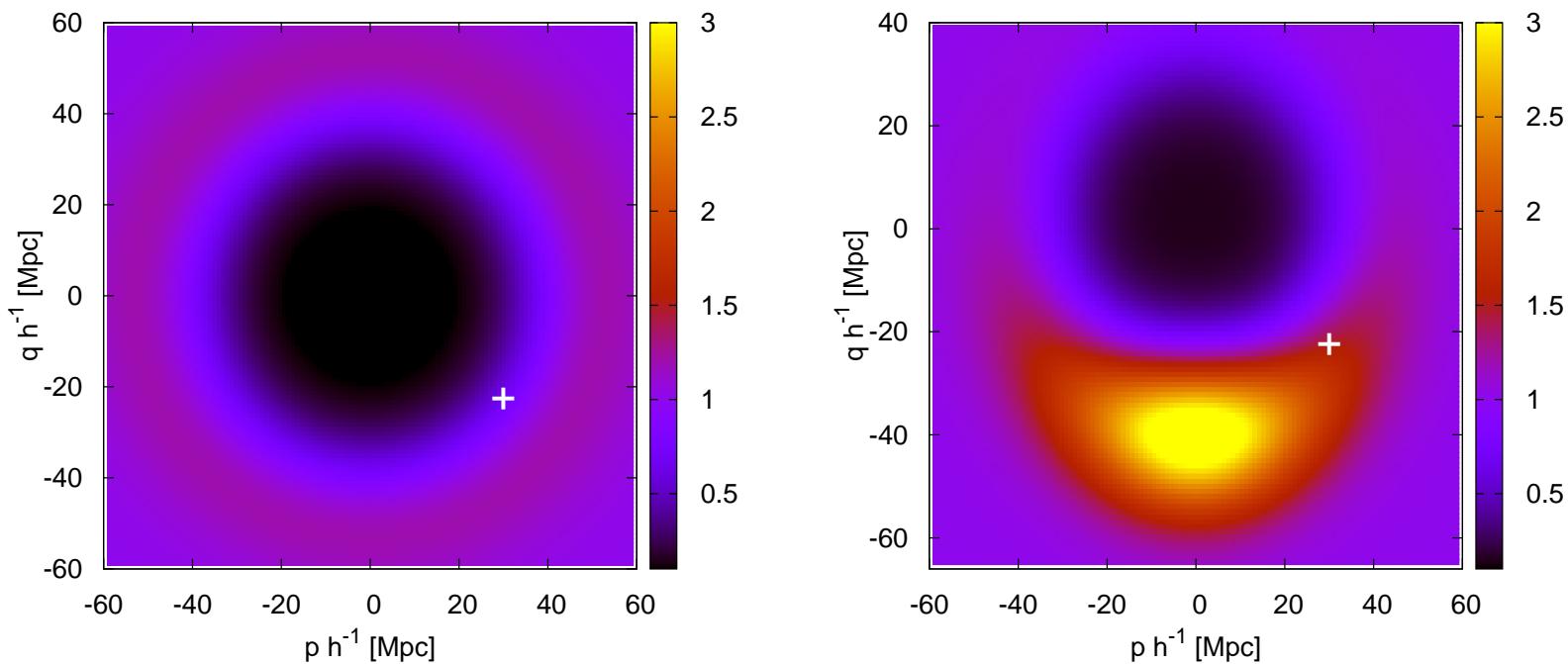


Simplify to  $P = Q = 0$ ;  $S = r \left( \frac{r}{1h^{-1}\text{Mpc}} \right)^{\alpha-1}$

$$\begin{aligned} M &= M_0 r^3 [1 + \delta_M(r)], \\ \delta_M(r) &= \frac{1}{2} \delta_0 \left( 1 - \tanh \frac{r - r_0}{2\Delta r} \right), \end{aligned}$$

$-1 \leq \delta_0 < 0$  underdensity at  $r \rightarrow 0$ ;  $\delta_M \rightarrow 0$  as  $r \rightarrow \infty$ .

# LTB and Szekeres profiles



- Fix  $\Delta r = 0.1r_0$ ,  $\varphi_{obs} = 0.5\pi$
- LTB parameters:  $\alpha = 0$ ,  $\delta_0 = -0.95$ ,  $r_0 = 45.5 h^{-1}$  Mpc;  $r_{obs} = 28 h^{-1}$  Mpc,  $\vartheta_{obs} = \text{any}$
- Szekeres parameters:  $\alpha = 0.86$ ,  $\delta_0 = -0.86$ ;  $r_{obs} = 38.5 h^{-1}$  Mpc;  $r_{obs} = 25 h^{-1}$  Mpc,  $\vartheta_{obs} = 0.705\pi$ .

# Szekeres model ray tracing constraints

- Require Planck satellite normalized FLRW model on scales  $r \gtrsim 100 h^{-1} \text{Mpc}$ ; i.e., spatially flat,  $\Omega_m = 0.315$  and  $H_0 = 67.3 \text{ km/s/Mpc}$
- CMB temperature has a maximum  $T_0 + \Delta T$ , where

$$\Delta T(\ell = 276.4^\circ, b = 29.3^\circ) = 5.77 \pm 0.36 \text{ mK},$$

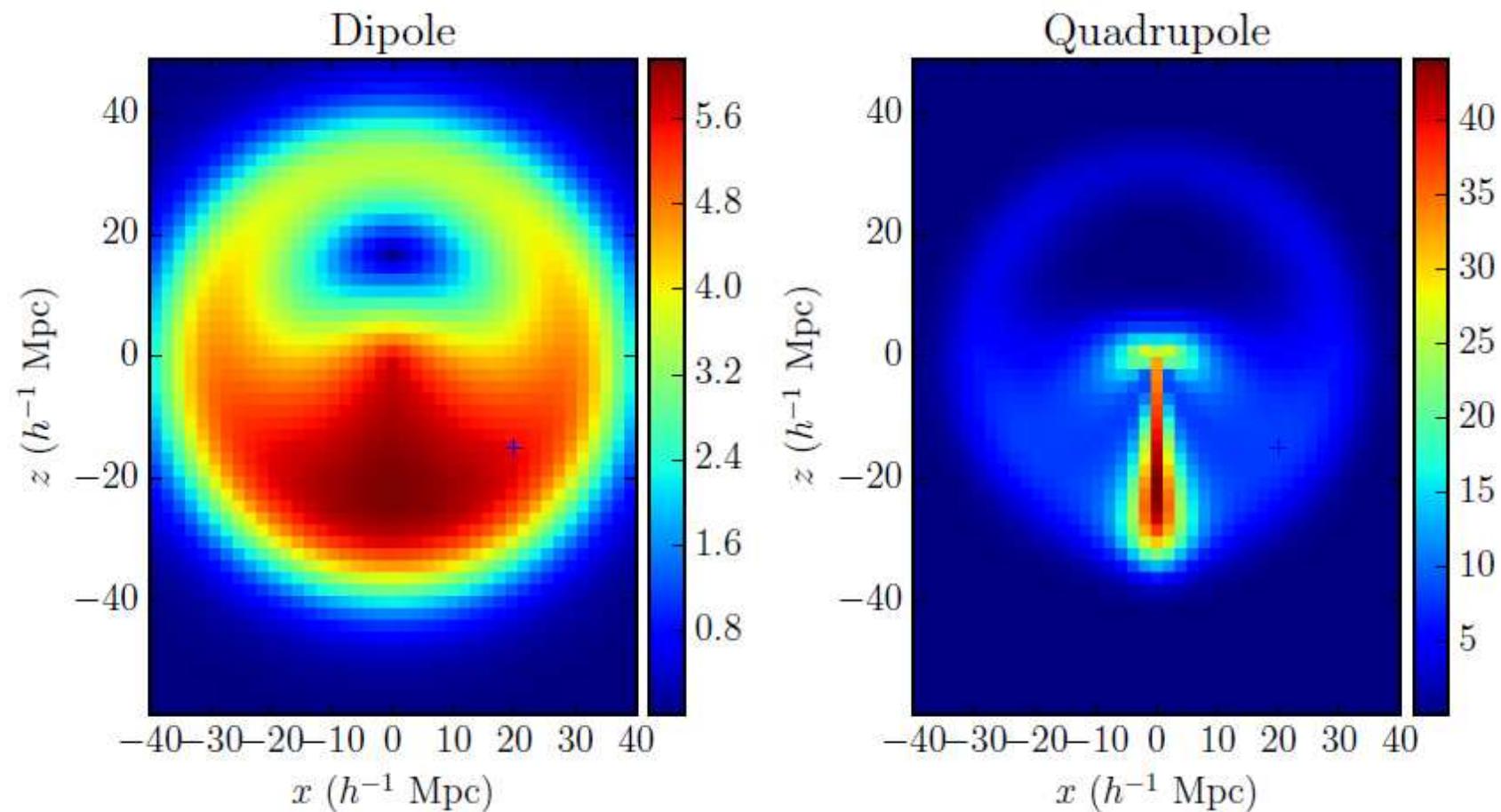
matching dipole amplitude, direction in LG frame

- CMB quadrupole anisotropy lower than observed

$$C_{2,CMB} < 242.2^{+563.6}_{-140.1} \mu\text{K}^2.$$

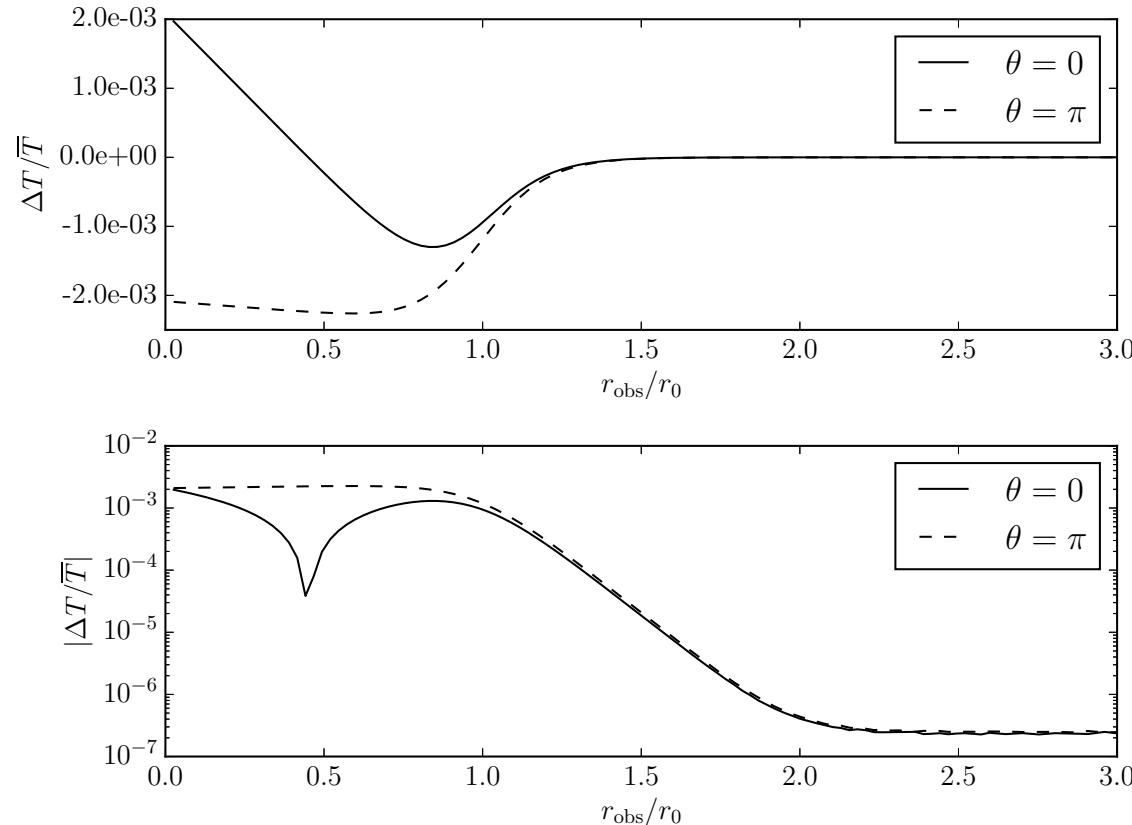
- Hubble expansion dipole (LG frame) matches COMPOSITE one at  $z \rightarrow 0$ , if possible up to  $z \sim 0.045$
- Match COMPOSITE quadrupole similarly, if possible

# CMB dipole, quadrupole examples



- Generate  $z_{\text{ls}}(\hat{\mathbf{n}})$  for each gridpoint
- $T = T_{\text{ls}}/(1 + z_{\text{ls}}); (T_{\text{max}} - T_{\text{min}})/2$  left (mK);  $C_2$  right ( $\mu\text{K}^2$ )

# Peculiar potential not Rees–Sciama



- Rees–Sciama (and ISW) consider photon starting and finishing from *average* point
- Across structure  $|\Delta T|/T \sim 2 \times 10^{-7}$
- Inside structure  $|\Delta T|/T \sim 2 \times 10^{-3}$

# Local expansion variation methodology

$$H_0(\ell, b, z) = \frac{\sum_i H_i w_{d,i} w_{z,i} w_{\theta,i}}{\sum_i w_{d,i} w_{z,i} w_{\theta,i}}, \quad \langle H_0 \rangle = \frac{1}{4\pi} \int d\Omega H_0(\ell, b, z)$$

$$\zeta_i = z_i + \frac{1}{2}(1 - q_0)z_i^2 - \frac{1}{6}(1 - q_0 - 3q_0^2 + j_0)z_i^3$$

$$H_i = c\zeta_i/d_i, \quad w_{d,i} = c\zeta_i d_i / (\Delta d_i)^2,$$

$$w_{z,i} = \frac{1}{\sqrt{2\pi}\sigma_z} \exp \left[ -\frac{1}{2} \left( \frac{z - z_i}{\sigma_z} \right)^2 \right],$$

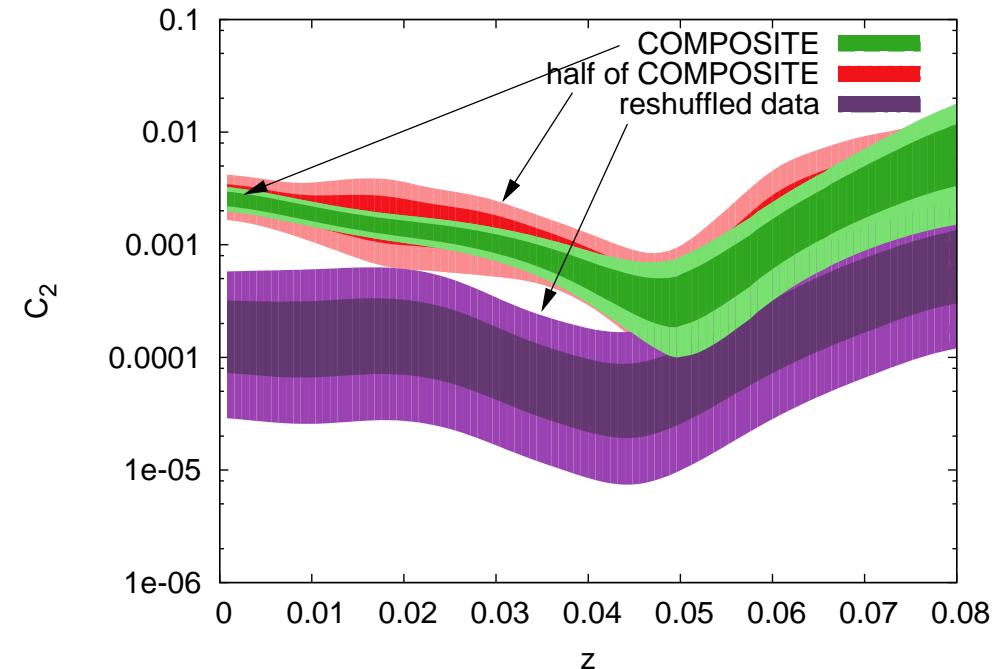
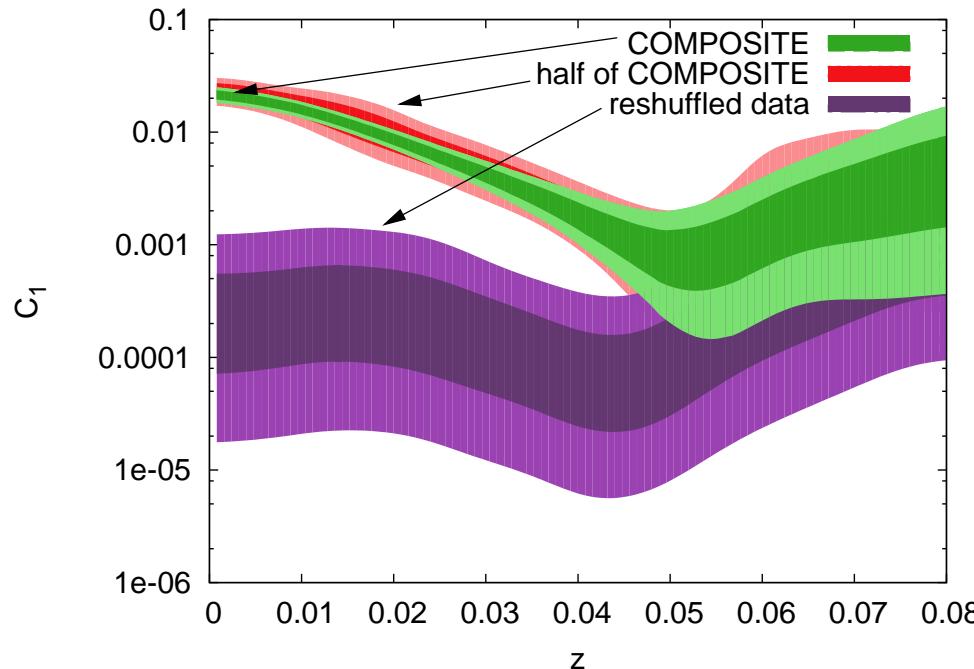
$$w_{\theta,i} = \frac{1}{\sqrt{2\pi}\sigma_\theta} \exp \left[ -\frac{1}{2} \left( \frac{\theta_i}{\sigma_\theta} \right)^2 \right],$$

$q_0 = -0.5275$ ,  $j_0 = 1$  ( $\Omega_m = 0.315$   $\Lambda$ CDM);  $\sigma_z = 0.01$ ,  
 $\sigma_\theta = 25^\circ$ ,  $\theta_i$  = angle between each source and boost apex.

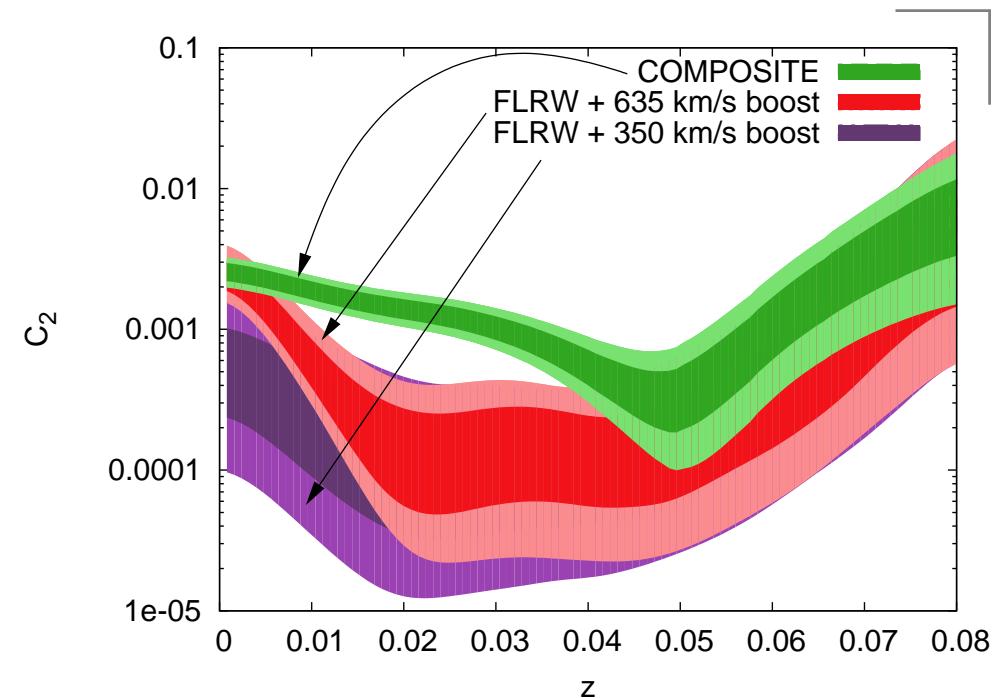
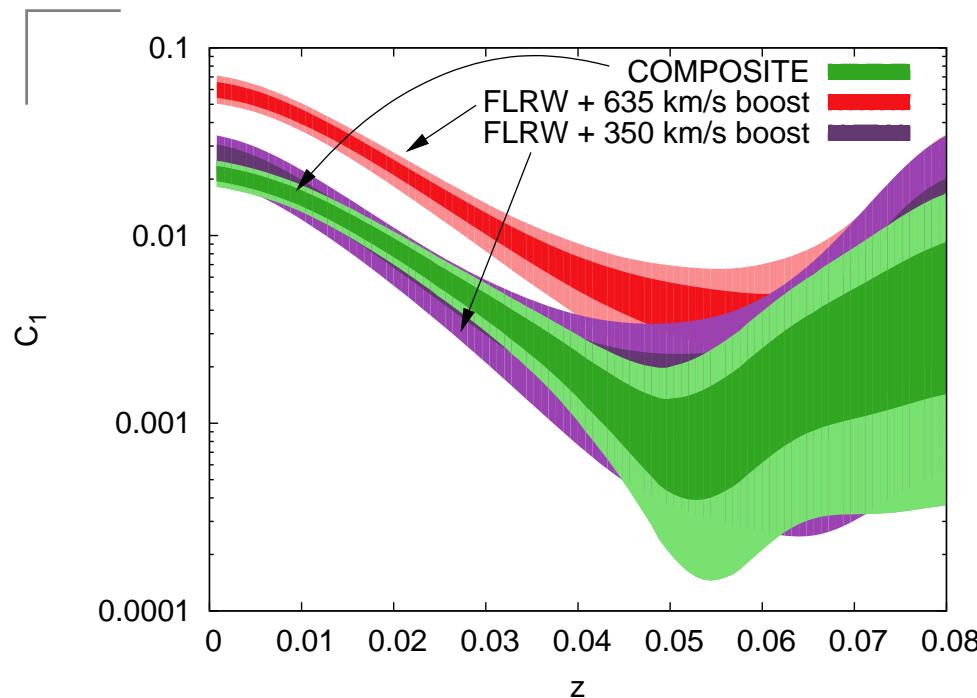
# Method and COMPOSITE data

$$\frac{\Delta H_0}{\langle H_0 \rangle} = \frac{H_0(l, b, z) - \langle H_0 \rangle}{\langle H_0 \rangle} = \sum_{l,m} a_{lm} Y_{lm},$$

Expand fractional Hubble expansion variation in multipoles,  
evaluate angular power spectrum:  $C_\ell = \frac{1}{2\ell+1} \sum_m |a_{\ell m}|^2$

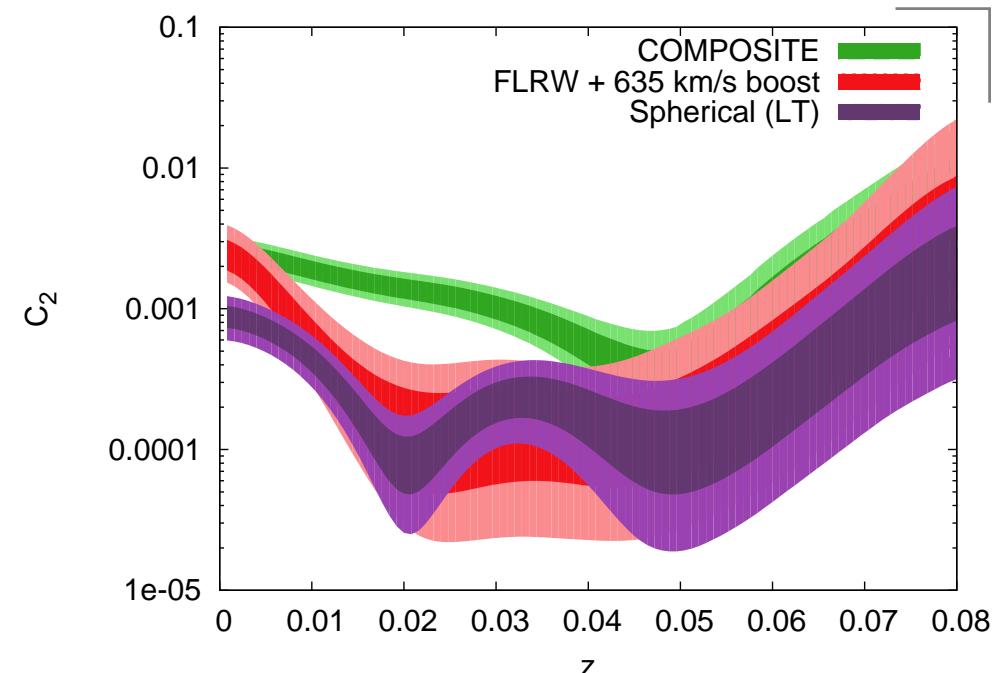
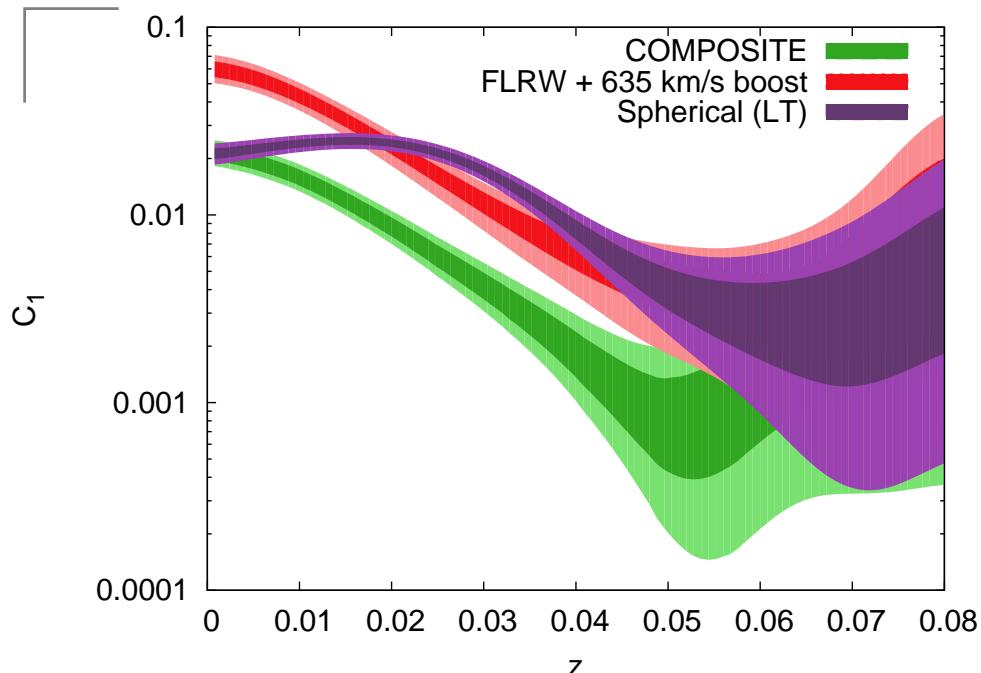


# FLRW model in CMB frame + LG boost



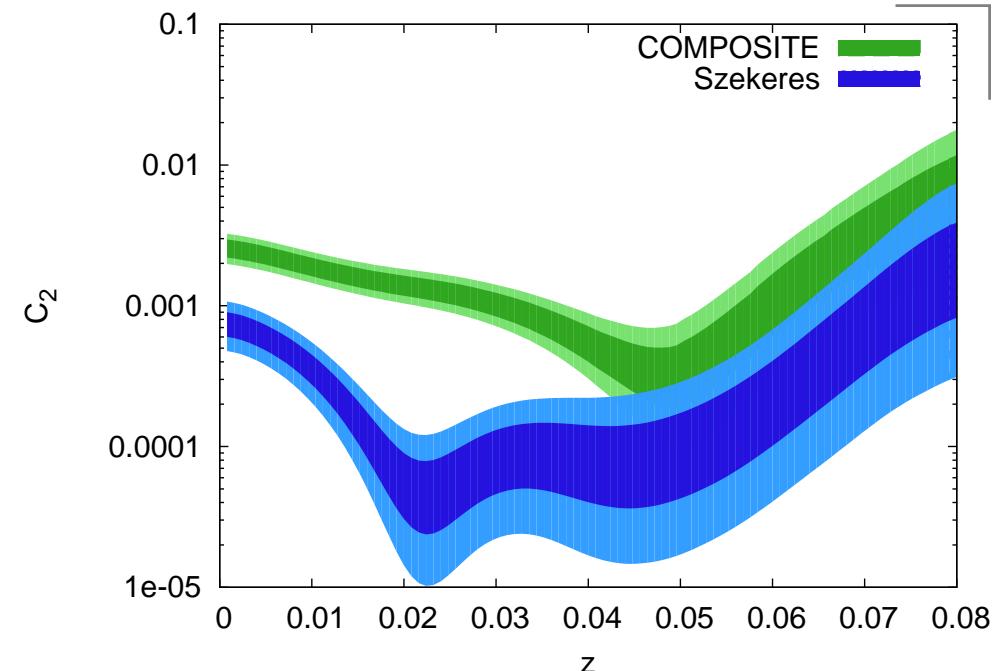
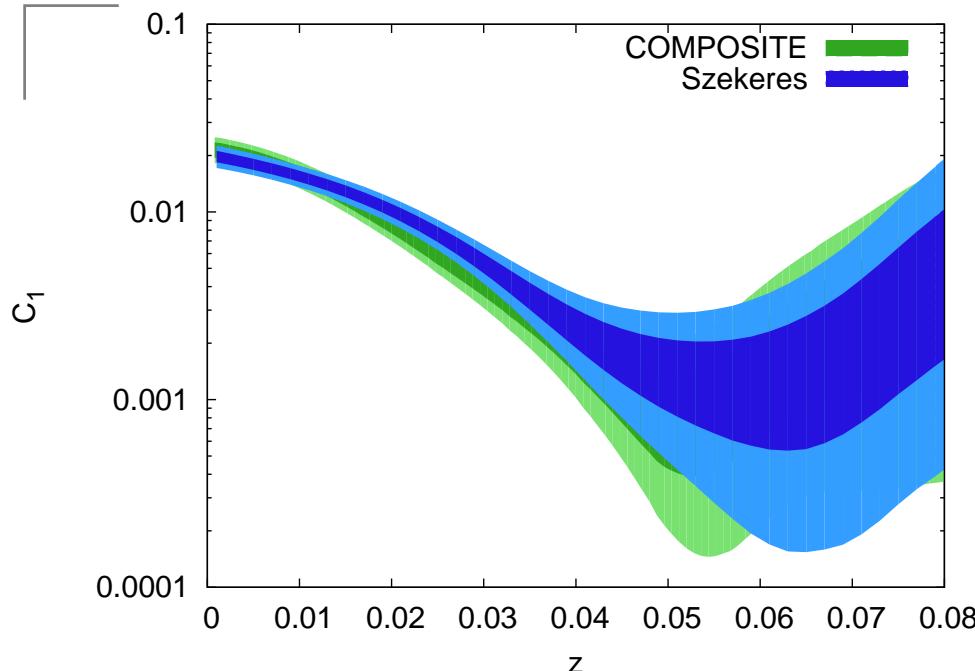
- Result of 100,000 mock COMPOSITE catalogues with same distance uncertainties
- For standard kinematic CMB dipole,  $H_0$  dipole too high over all  $z < 0.045$ ; quadrupole OK only at  $z \rightarrow 0$
- Dipole result: means bulk flow in standard approach

# LTB fit: $H$ dipole, quadrupole



- LTB dipole matches only at  $z \rightarrow 0$ , increases to close to FLRW plus boost case for larger  $z$
- Smaller insignificant quadrupole
- Differential expansion radially only; effective point symmetry on scale larger than inhomogeneity

# Best fit Szekeres: $H$ dipole, quadrupole

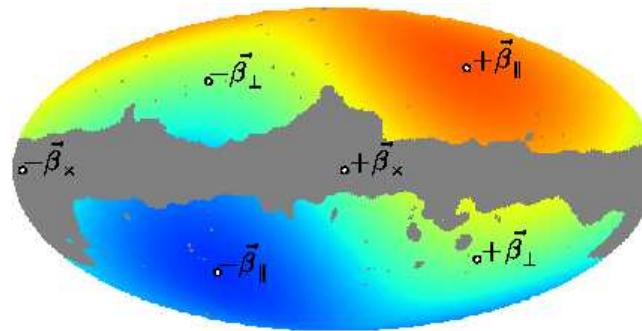


- Szekeres matches dipole on whole  $z < 0.045$  range
- Smaller insignificant quadrupole
- Note  $C_{2,CMB} = 8.26 \mu\text{K}^2$  30 times smaller than observed
- Possible additional smaller amplitude structures can add quadrupole (future work)

# Association with known structures

- Our galaxy is in a local void complex on a filamentary sheet (Tully et al 2008) joined to Virgo cluster.  
Dominant overdensity  $23 h^{-1}$ Mpc wide “Great Attractor”:
  - Near side Centaurus,  $z_{\text{LG}} = 0.0104 \pm 0.0001$ ,  
 $(\ell, b) = (302.4^\circ, 21.6^\circ)$
  - Far side Norma,  $z_{\text{LG}} = 0.0141 \pm 0.0002$ ,  
 $(\ell, b) = (325.3^\circ, -7.3^\circ)$
- Szekeres  $\delta\rho/\rho > 2$  ellipsoidal overdense region, spans  $0.003 \lesssim z_{\text{LG}} \lesssim 0.013$  (or  $16 h^{-1} \lesssim D_L \lesssim 53 h^{-1}$ Mpc) and angles  $220^\circ < \ell < 320^\circ$ ,  $-60^\circ < b < 40^\circ$
- Centaurus lies inside; Norma just outside
- Adding structures at larger distances (Perseus–Pisces) will change far side alignment

# Systematics for CMB



- Define nonkinematic foreground CMB anisotropies by

$$\Delta T_{\text{nk-hel}} = \frac{T_{\text{model}}}{\gamma_{\text{LG}}(1 - \boldsymbol{\beta}_{\text{LG}} \cdot \hat{\mathbf{n}}_{\text{hel}})} - \frac{T_0}{\gamma_{\text{CMB}}(1 - \boldsymbol{\beta}_{\text{CMB}} \cdot \hat{\mathbf{n}}_{\text{hel}})}$$
$$T_{\text{model}} = \frac{T_{\text{dec}}}{1 + z_{\text{model}}(\hat{\mathbf{n}}_{\text{LG}})}, \quad T_0 = \frac{T_{\text{dec}}}{1 + z_{\text{dec}}}$$

$z_{\text{model}}(\hat{\mathbf{n}}_{\text{LG}})$  = anisotropic Szekeres LG frame redshift;  
 $T_0$  = present mean CMB temperature

- Constrain  $\frac{T_{\text{model}}}{\gamma_{\text{LG}}(1 - \boldsymbol{\beta}_{\text{LG}} \cdot \hat{\mathbf{n}}_{\text{hel}})} - T_{\text{obs}}$  by Planck with sky mask

# Large angle CMB anomalies?

Anomalies (significance increased after Planck 2013):

- power asymmetry of northern/southern hemispheres
- alignment of the quadrupole and octupole etc;
- low quadrupole power;
- parity asymmetry; ...

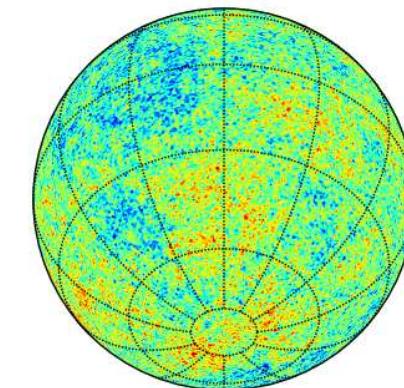
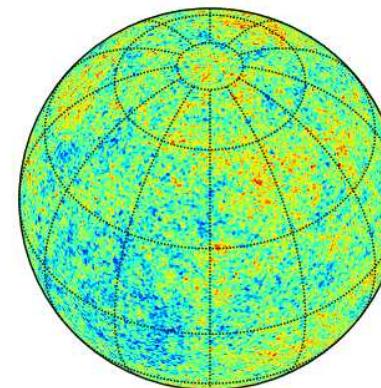
Critical re-examination required; e.g.

- light propagation through Hubble variance dipole foregrounds may differ subtly from Lorentz boost dipole
- dipole subtraction is an integral part of the map-making;  
is galaxy correctly cleaned?
- Freeman et al (2006): 1–2% change in dipole subtraction may resolve the power asymmetry anomaly.

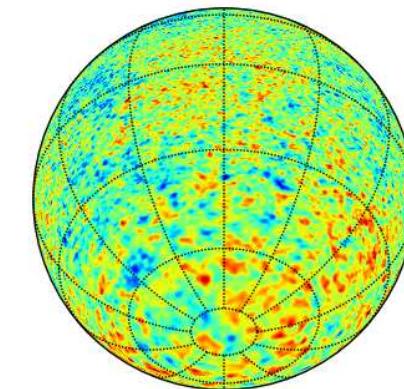
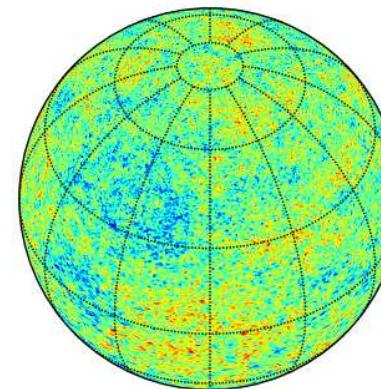
# Planck results arXiv:1303.5087

Boost dipole from  
second order effects

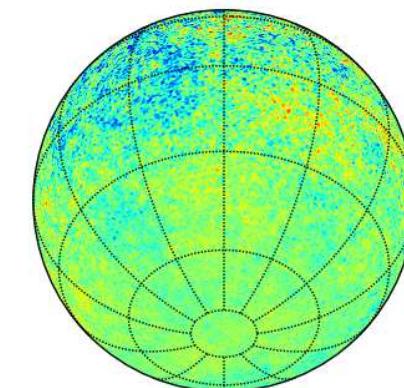
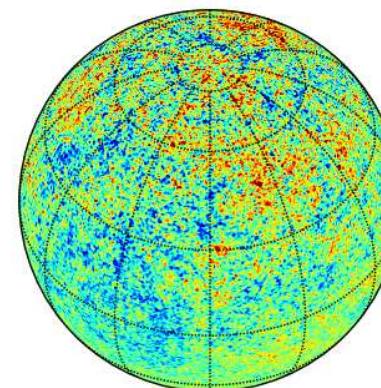
*Original*



*Aberration  
(Exaggerated)*

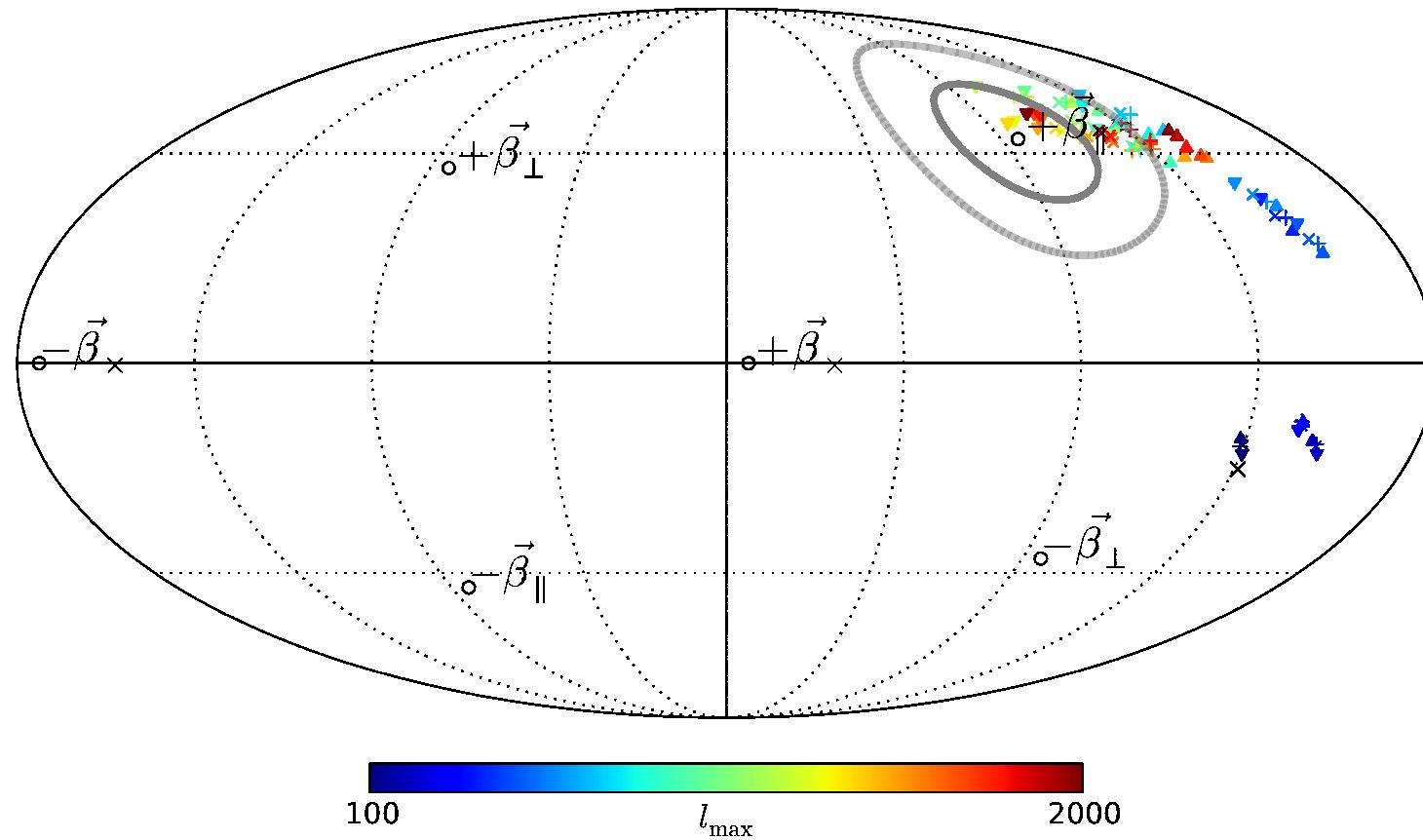


*Modulation  
(Exaggerated)*



Eppur si muove?

# Planck Doppler boosting 1303.5087



- Dipole direction consistent with CMB dipole  $(\ell, b) = (264^\circ, 48^\circ)$  for small angles,  $l_{\min} = 500 < l < l_{\max} = 2000$
- When  $l < l_{\max} = 100$ , shifts to WMAP power asymmetry modulation dipole  $(\ell, b) = (224^\circ, -22^\circ) \pm 24^\circ$

# Conclusion/Outlook

- Claim: GR below SHS can resolve bulk flow dilemmas
- Independently, kinematic dipole rejected at 99.5% confidence in radio galaxy number counts (Rubart & Schwarz 2013) – direction same
- A 0.5% nonkinematic anisotropy on  $\lesssim 65 h^{-1} \text{Mpc}$  scales has profound implications for cosmology
- Local value of  $H_0$  by Riess et al analysis higher than global value (Bolejko et al, to appear)
- CMB dipole subtraction *will be different* with modellable consequences for large angle anomalies
- Generalization for ensembles of realistic Szekeres models relevant for ISW amplitude calculation (needs assumptions about backreaction, averaging)